Fundamentals of Heat Transfer

RADIATION

INTRODUCTION

What is radiation?

- Radiation is the energy emitted by matter in the form electromagnetic waves or protons as a result of the changes in the electronic configuration of the atoms or molecules because of their temperature.
- All bodies at a temperature above absolute zero emit thermal radiation.
- Thermal Radiation occurs approximately in the range of 0.1 μm to 100 μm , including the visible and the near, middle and far IR (infrared) regions

INTRODUCTION

Electromagenetic wave spectrum

Thermal Radiation occurs approximately in the range of 0.1 μm to 100 μm , including the visible and the near, middle and far IR (infrared) regions

 10^{9} 10^{8} Radio and TV waves 10^{7} 10^{6} 10⁵ 10^{4} Microwaves 10^{3} 10^2 Thermal Infrared 10 radiation Visible 10-1 Ultraviolet 10^{-2} 10^{-3} X-rays 10-4 10-5 ү-гауѕ 10-6 10^{-7} 10-8 Cosmic

Electrical

power waves

λ, μm

 10^{10}

Radiation depends on:

- Surface temperature
- Surface Radiation properties
- Geometry
- Radiation properties of the medium

In general thermal radiation has both directional (*depends* on the angle relative to the surface) and spectral (*depends* on wave length) dependence. In this course surfaces are taken to be diffuse and consider only total emission

Black Body (Ideal Surface):

- Absorbs all incident radiation regardless of wavelength (λ) or angle of incidence and reflect zero radiation. It is a diffuse surface emitter (NO preferred direction for emission)
- Radiation emitted by a black body per unit time and per unit surface area was determined experimentally by Joseph Stefan in 1879 ——— Stefan's law

Radiation

- Radiation is the energy emitted by matter in the form of electromagnetic waves (or photons) as a result of the changes in the electronic configurations of the atoms or molecules.
- Unlike conduction and convection, the transfer of energy by radiation does not require the presence of an intervening medium.
- We focus on thermal radiation
- All bodies at a temperature above absolute zero emit thermal radiation.

The maximum rate of radiation that can be emitted from a surface at an absolute temperature *Ts* (in K or R) is given by the **Stefan-Boltzmann law** as:

$$\dot{Q}_{\text{emit}} = \varepsilon \sigma A_s T_s^4 \tag{W}$$

• For 2 bodies at temps:

$$\dot{Q}_{\rm rad} = \varepsilon \sigma A_s \left(T_s^4 - T_{\rm surr}^4 \right) \tag{W}$$

Laws of Thermal Radiation

Spectral black body emissive power

Radiant Energy emitted by a black body at T per unit area of surface per unit wavelength (Planck's Law):

$$E_{b_{\lambda}} = \frac{C_1 \lambda^{-5}}{e^{C_2/\lambda T} - 1} \dots (1)$$

where

 $C_1(=2\pi hc_0)=3.742 \times 10^8 \, W(\mu m)^4/m^2; C_2(=hc_0/k)=1.4389 \times 10^4 (\mu m.K);$ k=1.38065 \times 10^{-23} J/K is the Boltzmann constant; h=6.6256 \times 10^{-34} J.s is the universal Planck constant; $c_0=2.998 \times 10^8 \, m/s$ is the speed of light in vacuum

Equation (1) is valid for a surface in vacuum or gas.
 For other media, C₁ is replaced by C₁/n², where n is the index of refraction of the medium

CP 304 - Lecture # 5

Laws of Thermal Radiation

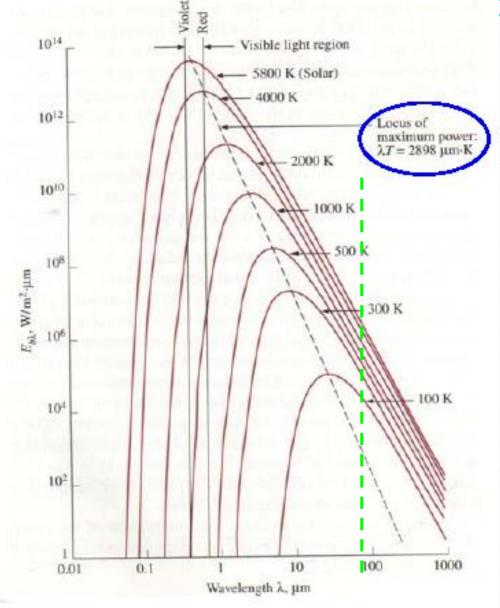


Figure 1: Emissive power vs λ at diff. temp

- 1. The emitted radiation is a continuous function of time. At any specific temperature, it increases with λ , reaches a peak, and then decreases with increasing λ
- 2. At any λ , the amount of emitted radiation increases with increasing temperature
- 3. As temperature increases , the curves shift to the left of the shorter λ
- 4. The radiation emitted by the sun which is considered as a black body at 5800K, reaches its peak at visible region.

Wien's displacement law

The wavelength corresponding to the peak radiant flux is inversely proportional to temperature

$$(\lambda T)_{\text{max power}} = 2897.8 \,\mu\text{m.K.}...(2)$$

If you can integrate spectral blackbody emissive power over the wavelength spectrum can give total black body emissive power

$$\mathbf{E}_{\mathbf{b}} = \int_{0}^{\infty} \mathbf{E}_{\mathbf{b}_{\lambda}} \mathbf{d}\lambda$$

$$E_{b_{\lambda}} = \frac{C_1 \lambda^{-5}}{e^{C_2/\lambda T} - 1}$$

$$E_{b_{\lambda}} = \frac{C_{1}\lambda^{-5}}{e^{C_{2}/\lambda T} - 1} \Rightarrow \frac{E_{b_{\lambda}}}{T^{5}} = \frac{C_{1}}{(\lambda T)^{5} \left[e^{C_{2}/\lambda T} - 1\right]}$$

$$\frac{E_b}{T^5} = C_1 \int_0^\infty \frac{d\lambda}{(\lambda T)^5 \left[e^{C_2/\lambda T} - 1\right]}$$

Let,
$$x = \frac{1}{\lambda T}, \Rightarrow dx = -\frac{1}{\lambda^2 T}$$

$$d\lambda = \frac{(\lambda T)^2}{T} dx \Rightarrow d\lambda = \frac{1}{x^2 T} dx$$

$$\frac{E_{b}}{T^{5}} = \frac{C_{1}}{T} \int_{0}^{\infty} \frac{x^{3} dx}{e^{C_{2}x} - 1}$$

$$\frac{E_b}{T^4} = C_1 \int_0^\infty x^3 \left[1 - e^{C^{2x}} \right]^{-1} dx$$

$$= C_1 \int_0^\infty x^3 \left[1 + e^{C_2 x} + 1 / 2e^{2C_2 x} + \dots \right] dx$$

Integrate by parts:

$$\frac{E_b}{T^4} = \frac{\sigma C_1}{C_2^4} \left[1 + \frac{1}{2^4} + \frac{1}{3^4} + \dots \right]$$

$$= \frac{\sigma C_1}{C_2^4} \frac{\pi^4}{90}$$

Laws of Thermal Radiation

The Stephan-Boltzmann Law

Thus, Integrating Eq. (1) over all wavelengths to give the heat transferred by radiation from a black body we get

$$E_b = \sigma T^4 \left(\frac{W}{m^2}\right) \dots (3)$$

where σ is Stephan-Boltzmann constant (=5.67× 10⁻⁸ W/m². K⁴), T is the absolute temperature of the surface in K

Laws of Thermal Radiation

For two black bodies at temperatures T1 and T2, the maximum relative interchange between them is given by

$$E_{b,1} = \sigma T_1^4 \dots (4a)$$

$$E_{b,2} = \sigma T_2^4 \dots (4b)$$

$$E_{b,1-2} = \sigma \left(T_1^4 - T_2^4\right) \dots (4c)$$

Radiation properties

Emissivity(ε)

Is the ratio of the radiation emitted by the surface at a given temp. to the radiation emitted by a blackbody at the same temp. It is a measure of how closely a surface approximates a black ($\epsilon = 1$)

$$0 \le \varepsilon \le 1 \tag{5}$$

Total hemispherical emissivity (or "average emissivity")

Is the ratio of radiation emitted surface to the radiation emitted by blackbody E(T)

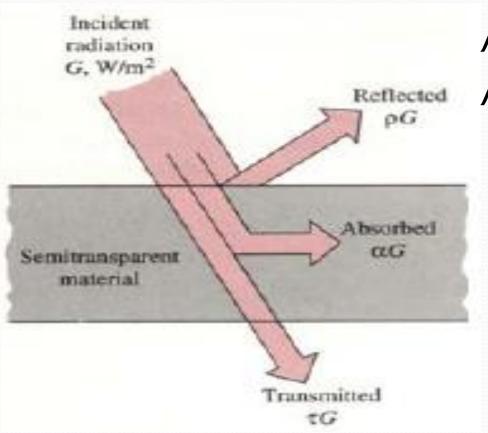
 $\varepsilon(T) = \frac{E(T)}{E_b(T)} \tag{6}$

Spectral directional emissivity

The ratio of intensity of radiation emitted by the surface at a specific wave length in specific direction to the intensity of radiation emitted by blackbody at the same temp. at the same wavelength.

Absorptivity (α), Reflectivity (ρ) & Transmissivity (τ)

Not all surfaces are ideal emitters and absorbers, they are not "Black" but "Grey"



$$\rho G + \varepsilon G + \tau G = G \quad (7)$$

$$\rho + \varepsilon + \tau = 1 \quad (8)$$

Most solids except glass and some plastics are opaque, thus $\tau = 0$

Fig. 6: Radiation incidence

Absorptivity vs. Emissivity

- Kirchhoff's law: emissivity and absorptivity are equal when <u>thermal equilibrium exists</u> (i.e., surface temperature = source temperature of incident radiation);
- In practice, the emissivity and absorptivity of gray surfaces are equal even when thermal equilibrium does not exist This is experimentally observed fact rather than a universal rule! This rule must not be applied to non-gray real surfaces!!

Absorptivity vs. Emissivity

- Gray and Diffuse assumption very commonly used in radiation calculations
- Radiation calculation is quite complex: one of the reasons is the wavelength and direction dependence of radiation properties, e.g.;
 - Diffuse surface: its radiation properties are independent of direction;
 - Gray surface: its radiation properties are independent of wavelength

$$\varepsilon(T) = \varepsilon_{\lambda}(T) = \varepsilon_{\hat{s}}(T) = \text{constant....}(9)$$

$$\alpha(T) = \alpha_{\hat{s}}(T) = \text{constant}....(10)$$

Absorptivity vs. Emissivity

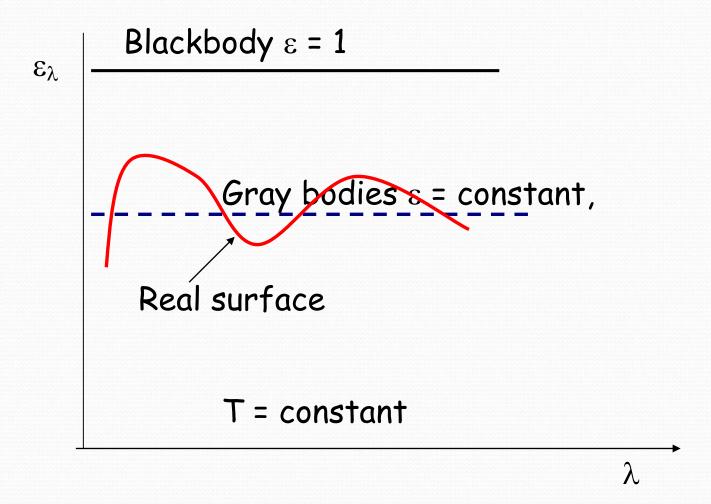
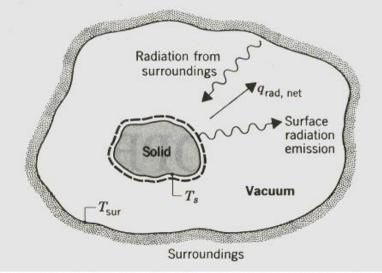


Fig. 7: Variation of Emissivity with wavelength

Kirchhoff's Law

 When two bodies A and B of areas A1 and A2 are in a large isothermal enclosure:



Energy absorbed by A from the enclosure

$$A_1\alpha_1I_1....(11)$$

Energy given out by A:

$$E_1I_1$$
....(12)

Kirchhoff's Law

 At equilibrium these quantities will be equal. In the same way the energy emitted by B will be equal to the energy received

$$A_1 \cdot \alpha_1 \cdot I_1 = A_1 \cdot E_1 \cdot \dots (13)$$

and

$$A_2 \cdot \alpha_2 \cdot I_2 = A_2 \cdot E_2 \dots (14)$$

• But $I_1 = I_2$, giving

$$I = \frac{E_1}{\alpha_1} = \frac{E_2}{\alpha_2} = \frac{E}{\alpha}$$
....(15)

 Equation (15) is general for any body and is termed Kirchhoff's law

Kirchhoff's Law

- Emissivity of a body is defined as the ratio of its emissive power to that of a black body Eq. (6)
- Since E/α is constant for all bodies, then

$$\frac{E}{\alpha} = \frac{E_b}{\alpha_b} \dots (16)$$

and

$$\varepsilon = \frac{E}{E_b} = \frac{\alpha}{\alpha_b} = \frac{\alpha}{1} = \alpha \dots (17)$$

Heat Transfer by Radiation

- A body of emissivity ε , at an absolute temperature T_1 emits energy $\varepsilon \sigma T_1^4$ per unit area.
- If the surroundings are black, they reflect back none of this radiation, but if they are at an absolute temperature T_2 they will emit radiation σT_2^4
- If the body is grey it will absorb a fraction ε, so that the net radiation per unit area from the grey body will be

$$\dot{q} = \varepsilon \cdot \sigma \cdot \left(T_1^4 - T_2^4\right) \dots (18)$$

Heat Transfer by Radiation

- For a material that does not behave as a grey body but as a **reflective emitter**, the absorptivity of the surface at T₁ for radiation from surroundings at T₂ will be αT₂.
- This will not be equal to its emissivity at T_1 (εT_1) but to its emissivity at T_2 (εT_2).
- Under these conditions the general equation for the net exchange of heat becomes:

$$\dot{q} = \sigma \cdot \left(\varepsilon_{T_1} \cdot T_1^4 - \varepsilon_{T_2} \cdot T_2^4\right) \dots (19)$$